

**FREQUENCY COMPARISON (H\_MASER 140 0810) - (LNE-SYRTE-FO1)  
For the period MJD 58419 to MJD 58449**

The primary frequency standard LNE-SYRTE-FO1 has been compared to the hydrogen maser 140 0810 of the laboratory, during a measurement campaign between MJD 58419 and 58449 (28<sup>th</sup> October 2018 – 27<sup>th</sup> November 2018). The fountain operation covered 93.3% of the estimation period.

The mean frequency difference at the middle date of the period is given in the following table:

Period (MJD)	Date of the estimation	$y(\text{HMaser}_{140\ 0810} - \text{FO1})$	$u_B$	$u_A$	$u_{\text{link} / \text{maser}}$
58419 – 58449	58434	-6463.7	3.2	2.5	0.5

*Table 1: Results of the comparison in  $1 \times 10^{-16}$ .*

The FO1 fountain was operated in the same mode during all the period: the interrogating signal synthesis is based on the multiplication of a 100 MHz signal provided by a cryogenic oscillator phase locked to the maser 140 0810. It uses a synthesizer to lock the microwave signal to the atomic resonance. The frequency difference between the maser and the fountain is deduced from the average correction applied to the synthesizer.

**Average value and statistical uncertainty**

The frequency data are averaged over 0.2 day intervals. We then perform a linear unweighted fit to the average data points to determine the average frequency at the middle date of the interval, as given in Table 1. The statistical uncertainty  $u_A$  is estimated using the Allan variance of the frequency residuals, after removing the drift. We estimate a conservative statistical uncertainty  $u_A$  of  $2.5 \times 10^{-16}$ .

We verified the result by applying a second method. We calculated the accumulated phase by integrating the data points, assuming a constant frequency during each segment, and during the dead times of the fountain operation. The average frequency is then obtained by dividing the total accumulated phase by the calibration period duration. The processing has been performed with segments of 0.01, 0.1 and 1 day duration. The results are in agreement with the values given in Table 1 within  $0.2 \times 10^{-16}$ , which is consistent with the estimation of the statistical uncertainty  $u_A$  and the uncertainty due to the link.

## Accuracy

The frequency is corrected from the quadratic Zeeman, the black body radiation, the cold collision (+ cavity pulling), the distributed cavity phase shift and the microwave lensing shifts, and at last the redshift. The cold collision (+ cavity pulling) correction is based on alternating measurements at full density for 50 cycles and at half density for 100 cycles, by changing the state selection frequency and readjusting the microwave power to keep the selection transition probability at maximum. The uncertainty in this correction is mostly statistical. The following table summarizes the budget of the systematic corrections and their associated uncertainties. The accuracy is the quadratic sum of all the systematic uncertainties.

	Correction ( $10^{-16}$ )	Uncertainty ( $10^{-16}$ )
Quadratic Zeeman effect	-1281.00	0.40
Black body radiation	169.50	0.60
Cold collisions + cavity pulling	124.19	1.44
Distributed cavity phase shift	-0.07(*)	2.40
Microwave lensing	-0.65	0.65
Microwave spectral purity&leakage	0	1.00
Ramsey & Rabi pulling	0	0.20
Second order Doppler effect	0	0.10
Background gas collisions	0	0.30
<b>Total</b>	<b>-988.03</b>	<b>3.16</b>
Redshift	- 69.08	0.25
<b>Total with redshift</b>	<b>-1057.11</b>	<b>3.17</b>

Table 2: Budget of systematic corrections and uncertainties for SYRTE-FO1 fountain for the MJD 58419 – 58449 period

$$u_B = 3.2 \times 10^{-16}$$

(\*) From now on, we apply an updated correction for the DCP shift: as detailed in the paper J. Guéna et al. IEEE TUFFC, 59, 3, 2012, this correction accounts for an effective tilt of the fountain launch direction along the microwave feed axis. From recent asymmetric feeding measurements, the effective tilt is compatible with zero, therefore the previous correction for the  $m=1$  term of the DCP is no more justified. Only remains the  $m=0$  term of the correction,  $-7 \times 10^{-18}$  as given in Table 2.

## Uncertainty of the link

The uncertainty of the link is the quadratic sum of 2 terms:

-A possible effect of phase fluctuations introduced by the cables that connect the primary standard to the maser. A new characterization of the signal distribution leads to a still conservative value of  $0.5 \times 10^{-16}$ .

-The uncertainty due to the dead times of the frequency comparison.

We have updated the estimation of this contribution, applying the method described in *Metrologia*, vol. 44, pp 91-96, 2007, as we did for the initial calibration reports of the LNE-SYRTE Strontium SFS. The maser noise model includes a white frequency noise component of  $5 \times 10^{-16}$  at 1 d and a flicker frequency noise component of  $5 \times 10^{-16}$  at 1 d, which is pessimistic especially for short averaging periods. We applied the method to the dead times longer than 600 s and obtained a stability degradation of  $0.2 \times 10^{-16}$ .

