

## Evaluation of PTB primary caesium fountain frequency standard CSF2 between MJD 57114 - MJD 57129

PTB's primary caesium fountain frequency standard CSF2 was operated between MJD 57114, 0:00 UTC and MJD 57129, 0:00 UTC. Frequency comparisons were made with respect to PTB hydrogen maser H9, BIPM code 1400509.

The relative frequency instability of CSF2 was  $2.16 \times 10^{-13} (\tau/s)^{-1/2}$  during the 15 days. The actual measurement time amounts to 98.51% of the  $15 \times 24$  hours. This results in a statistical uncertainty  $u_A = 0.19 \times 10^{-15}$  assuming that white frequency noise is the dominant noise source.

For the uncertainty due to the clock link  $u_{Lab} = 0.02 \times 10^{-15}$  is obtained by taking into account the actual measurement time. Finally, the estimated uncertainty for the link to TAI for 15 days is  $u_{TAI} = 0.12 \times 10^{-15}$ .

Frequency corrections for the following effects were applied to the raw data:

- Zeeman effect (magnetic field along the atoms' trajectory)
- blackbody effect (thermal radiation along the atoms' trajectory)
- gravitational red-shift and relativistic Doppler effect
- cold collisions effect
- cavity phase effect
- microwave lensing effect

The CSF2 standard uncertainty  $u_B$  is estimated as  $0.33 \times 10^{-15}$  ( $1 \sigma$ ) for the relevant period.

### **Table of results of CSF2 compared to hydrogen maser H9 (1400509)**

Interval of evaluation	MJD 57114, 0:00 UTC - MJD 57129, 0:00 UTC
Fractional dead time	1.49 %
Resulting frequency difference	$y(\text{CSF2} - \text{H9}) = 74.33 \times 10^{-15}$
Type A uncertainty $u_A$ ( $1 \sigma$ )	$0.19 \times 10^{-15}$
Type B uncertainty $u_B$ ( $1 \sigma$ )	$0.33 \times 10^{-15}$
Link to clock $u_{Lab}$ ( $1 \sigma$ )	$0.02 \times 10^{-15}$
Link to TAI $u_{TAI}$ ( $1 \sigma$ )	$0.12 \times 10^{-15}$ (15 days)
Combined uncertainty ( $1 \sigma$ )	$0.40 \times 10^{-15}$

## **Operation mode of CSF2**

Formerly CSF2 was operated in an autonomous mode by steering a quartz oscillator to the atomic resonance frequency and measuring the frequency difference between the quartz oscillator and the hydrogen maser with a commercial phase comparator [1]. Since 2011 CSF2 is operated in a non-autonomous mode: As in other fountains, for this purpose the quartz oscillator is locked to the hydrogen maser. The microwave signal for the fountain is synthesized based on the quartz oscillator and a DDS synthesizer locked to the quartz oscillator. The output frequency of the DDS synthesizer is digitally steered to the atomic resonance frequency by evaluating the caesium transition probability measured by the fountain. The frequency difference between CSF2 and the hydrogen maser is thus obtained from the monitored synthesizer frequency settings.

## **Type A (statistical) uncertainty of CSF2**

For the TAI scale unit measurement at hand the atoms were loaded from the background gas into the molasses. The frequency instability was measured to be  $2.16 \times 10^{-13} (\tau/s)^{-1/2}$ .

In 2010, a new microwave frequency synthesis setup [2] identical to the one used in the fountain PTB-CSF1 has been introduced in the CSF2 electronics setup. Because it had been demonstrated that the new synthesis setup is capable of providing instabilities below the  $10^{-16}$  level, the statistical uncertainty of CSF2 frequency measurements is no more limited at the  $7 \times 10^{-16}$ -level as before [1].

For this reason the statistical uncertainty of the current TAI scale unit measurement was calculated with the assumption of white frequency noise for the total measurement interval.

## **Type B (systematic) uncertainty of CSF2**

Detailed descriptions of the systematic uncertainty contributions of the PTB fountain CSF2 have been published elsewhere [1], [3]. Here we only report some details about the current methods for evaluating the quadratic Zeeman shift and the collisional shift because they differ from our previously employed and described methods. We also briefly comment on the reduced uncertainty due to the electronics.

The average value of the quadratic Zeeman shift is determined by automated periodic measurements of the frequency of the  $(F=3, m=-1) \rightarrow (F=4, m=-1)$  transition during the relevant period. Therefore the uncertainty of the quadratic Zeeman correction is now dominated by the statistical measurement uncertainty and amounts to  $0.010 \times 10^{-15}$ .

To prepare the fountain for future evaluations at increased atomic density, a new method to measure the collisional shift has been implemented using rapid adiabatic passage [4,5]. The required microwave pulses are applied in the state selection cavity. By switching from the full microwave pulse to a pulse that is cut off at the exact pulse center, it is possible to reduce the density of the atomic cloud to 50% of its original value at any position in the atomic cloud, leaving the relative distribution unchanged.

During the present evaluation, the fountain was alternately operated at high (300 shots) and low density (400 shots) using the rapid adiabatic passage method for the density variation. The collisional shift was thus evaluated online during the fountain evaluation while the differential measurement eliminates the effect of frequency drifts of the hydrogen maser reference. As in previous evaluations, the result of this evaluation is a slope factor which gives – multiplied with the actual number of atoms – the collisional frequency shift correction [1].

For the present evaluation we calculate the collisional shift based on the measured relative atom numbers during the present evaluation and a slope factor, which is the weighted average of slope factors obtained during the present evaluation and during the two previous evaluations. The collisional shift uncertainty is given by the statistical measurement uncertainty of the slope factors. In total we obtain a collisional shift uncertainty of  $0.25 \times 10^{-15}$ .

The new microwave frequency synthesis setup [2] provides a better suppression of sidebands compared to the previously employed synthesis. The dominating 50 Hz sidebands are 65 dB below the carrier with an asymmetry of much less than 10%. Therefore the current conservative estimate of the uncertainty due to the electronics amounts to  $0.1 \times 10^{-15}$ .

As a result of a reevaluation the frequency correction due to microwave lensing [3], [6] has been slightly changed (within the uncertainty). Because we take half of this correction as the uncertainty, also the latter has been slightly changed.

Rabi and Ramsey pulling have been recently reevaluated theoretically and experimentally using more elaborate methods than before [7]. The resulting combined uncertainty is  $0.0013 \times 10^{-15}$ .

In the table below we report the type B uncertainty evaluation results valid for the evaluation at hand.

**Frequency shifts, corrections and type B uncertainties of CSF2 (parts in  $10^{15}$ ):**

Frequency shift	Correction	Uncertainty
Quadratic Zeeman shift	-100.247	0.010
Blackbody radiation shift	16.557	0.057
Gravity+relativistic Doppler effect	-8.567	0.006
Collisional shift	0.63	0.25
Cavity phase shift	-0.044	0.133
Microwave lensing	-0.072	0.036
AC Stark shift (light shift)		0.001
Majorana transitions		0.0001
Rabi and Ramsey pulling		0.0013
Electronics		0.10
Microwave leakage		0.10
Background gas collisions		0.05
Total type B uncertainty		0.33

## References

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